

# LiCoO<sub>2</sub> and LiCo<sub>1-x</sub>Al<sub>x</sub>O<sub>2</sub> thin film cathodes grown by pulsed laser ablation

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## Abstract

LiCoO<sub>2</sub> and LiCo<sub>0.5</sub>Al<sub>0.5</sub>O<sub>2</sub> thin films have been grown by pulsed laser ablation on SnO<sub>2</sub>-coated glass substrates. For both stoichiometries, the resultant films are dense and uniaxially textured films with the Li and Co layers parallel to the substrate. In general, to grow LiCo<sub>0.5</sub>Al<sub>0.5</sub>O<sub>2</sub> films, a laser flux roughly 80 mJ pulse<sup>-1</sup> higher than that used for LiCoO<sub>2</sub> films is required to achieve a similar deposition rate. LiCoO<sub>2</sub> films grown at  $T_s = 600^\circ\text{C}$  and  $p[\text{O}_2] = 2000$  mTorr have a typical grain size of  $\sim 100$  nm. For constant current cycling between 3.8 and 4.2 V at 5  $\mu\text{A}$ , the LiCoO<sub>2</sub> films have an initial discharge capacity of  $\sim 0.33$  Li per LiCoO<sub>2</sub> (89 (mA h) g<sup>-1</sup>) decreasing to  $\sim 0.18$  Li/LiCoO<sub>2</sub> (49 (mA h) g<sup>-1</sup>) after 100 cycles and have a continued capacity loss of  $\sim 0.25\%$  per cycle. The LiCo<sub>0.5</sub>Al<sub>0.5</sub>O<sub>2</sub> films grown to date have roughly 3 times less capacity than the LiCoO<sub>2</sub> films and apparently a large asymmetry between Li extraction and reintercalation. © 1999 Elsevier Science S.A. All rights reserved.

**Keywords:** LiCoO<sub>2</sub>; LiCo<sub>0.5</sub>Al<sub>0.5</sub>O<sub>2</sub>; Thin film cathodes; Li-ion batteries

## 1. Introduction

LiCoO<sub>2</sub> is one of the most promising cathode materials for rechargeable Li-ion batteries and is already employed in a large number of commercial Li-ion secondary batteries [1–3]. Typically, LiCoO<sub>2</sub> batteries are cycled between the discharged fully-lithiated state LiCoO<sub>2</sub> ( $\sim 3.8$  V vs. Li) and a charged half-delithiated state Li<sub>0.5</sub>CoO<sub>2</sub> ( $\sim 4.2$  V vs. Li) yielding a useable capacity of  $\sim 0.5$  Li per transition metal ion (136 (mA h) g<sup>-1</sup>). In comparison,  $\omega$ -phase Li<sub>3</sub>V<sub>2</sub>O<sub>5</sub>, which can be synthesized electrochemically from V<sub>2</sub>O<sub>5</sub>, has a potential capacity of 3 Li per V<sub>2</sub>O<sub>5</sub> or 1.5 Li per transition metal ion [4]. Additional Li can be extracted from Li<sub>0.5</sub>CoO<sub>2</sub> by increasing the charging voltage limit. In particular, Amatucci et al. find that LiCoO<sub>2</sub> can be fully delithiated to form metastable CoO<sub>2</sub> upon charging to 5.2 V [5]. Understanding both the mechanism for the cycling-induced capacity loss in Li<sub>1-x</sub>CoO<sub>2</sub> over

the restricted intercalation range  $0 < x < 0.5$  and the potential for an increased capacity upon charging to  $x > 0.5$  requires a better understanding of the intrinsic electrochemical properties of LiCoO<sub>2</sub> as well as the effects of crystallinity, grain size and morphology on the electrochemical properties of real electrodes [6–9]. In addition, thin film growth of lithium-transition-metal oxides is central to thin film and micro-battery technology development.

Increasing the cell potential is another way to improve the stored energy density, provided that the charge capacity is not degraded. Recent calculations predict that chemical substitution of aluminum for cobalt in LiCo<sub>1-x</sub>Al<sub>x</sub>O<sub>2</sub> should increase the battery voltage [10]. As aluminum is both lighter than and less expensive than cobalt, this raises the tantalizing possibility that aluminum substituted LiCo<sub>1-x</sub>Al<sub>x</sub>O<sub>2</sub> could be a cheaper, lighter, higher voltage cathode material than the current standard, LiCoO<sub>2</sub>. Initial experiments, on bulk samples, do find an increased cathode voltage upon Al substitution albeit with a substantially reduced capacity [10].

In this work, we report on the pulsed laser deposition [11,12] and electrochemical characterization of LiCoO<sub>2</sub> and LiCo<sub>0.5</sub>Al<sub>0.5</sub>O<sub>2</sub> thin films grown on SnO<sub>2</sub>-coated glass

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substrates. For both the  $\text{LiCoO}_2$  and  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$ , dense uniaxially textured films with the  $R\bar{3}m$  symmetry of the layered  $\alpha\text{-NaFeO}_2$  crystal structure were grown [13,14]. For both  $\text{LiCoO}_2$  and  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  films, the grain size increases with increased substrate temperature in the range 400–700°C with roughly 100 nm grains for  $\text{LiCoO}_2$  films grown at 600°C and  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  grown at 700°C. In general, the Al-substituted material must be grown roughly 100°C hotter than the stoichiometric  $\text{LiCoO}_2$  to obtain similar grain size and crystallinity. For the  $\text{LiCoO}_2$  films grown here, a  $\text{Co}_3\text{O}_4$  impurity phase has been identified via Raman scattering measurements but is not observed in the X-ray diffraction measurements. This  $\text{Co}_3\text{O}_4$  impurity phase is found to decrease with increased substrate temperature and increased oxygen partial pressure.

For constant current cycling between 3.8 and 4.2 V at 5  $\mu\text{A}$ , the  $\text{LiCoO}_2$  films have an initial discharge capacity of  $\sim 0.33$  Li per  $\text{LiCoO}_2$  (89 (mA h)  $\text{g}^{-1}$ ) which decreases to  $\sim 0.18$  Li/ $\text{LiCoO}_2$  (49 (mA h)  $\text{g}^{-1}$ ) after 100 cycles. After 100 cycles the capacity loss rate is about 0.25% per cycle. Increasing the charging limit from 4.2 to 4.6 V increases the discharge capacity by 50%. For the  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  films, the amount of Li which can be extracted on the first charge cycle is similar to that for  $\text{LiCoO}_2$ . However, the subsequent capacity of the films grown to date is substantially lower, only  $\sim 0.1$  Li/ $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  (33 (mA h)  $\text{g}^{-1}$ ). An initial capacity loss rate of 6% per cycle decreases to 1% per cycle after 50 cycles.

## 2. Experimental procedures

All the films grown in this study were grown by pulsed laser ablation from nominally phase-pure stoichiometric ceramic  $\text{LiCoO}_2$  or  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  targets in a controlled atmosphere vacuum chamber using a 248 nm, 325–415 mJ pulse $^{-1}$ , excimer laser operating at 10 Hz. Typically, deposition runs were 20,000 pulses long. The target and heater are parallel, on the same axis, and roughly 8.5 cm apart. During deposition, the target is rotated at 1 to 10 rpm to ensure uniform wear from the ablating beam, which strikes the target at 45°. The focused spot size is  $1 \times 3$  mm at the target. To ensure good thermal contact, the  $\text{SnO}_2$ -coated 7059 glass substrates are mounted to the heater using silver paint. Crystalline  $\text{LiCoO}_2$  films have been grown over the substrate temperature ( $T_s$ ) range 300°C <  $T_s$  < 700°C and in the oxygen pressure ( $p[\text{O}_2]$ ) range 50 mT <  $p[\text{O}_2]$  < 2000 mTorr.

Average thickness measurements were made with a Dektak III profilometer. Corners of the  $\text{SnO}_2/7059$  substrates were masked during growth with  $\text{TiO}_2$  which is easily removed to allow these measurements. The metals stoichiometry was determined using inductively coupled plasma (ICP) analysis. For ICP, the  $\text{LiCoO}_2$  samples were initially dissolved in 1:1 HCl:H $_2$ O solution and then di-

luted to 10 vol% HCl. The phase, crystallinity and orientation of the films were determined using a four-circle Scintag diffractometer with a Cu-K $\alpha$  source. The Raman scattering measurements were performed in a 180° back-scattering configuration with an Ar-ion laser operating at 514.5 nm.

Electrochemical cycling experiments were conducted in an Ar-filled dry box using an Arbin Instruments four conductor battery test system with both the positive current and voltage leads connected to the cathode to mimic a standard three terminal electrochemical cell. For all tests, Li metal foil was used as both the anode and the reference voltage electrode. The electrolyte solution was either 1 M  $\text{LiClO}_4$  in propylene carbonate (PC) or 1 M  $\text{LiPF}_6$  in a 1:1 wt.% ethylene carbonate (EC)-dimethyl carbonate (DMC) solution. All measurements were made under constant current conditions with various voltage ranges. The specific charge capacities reported here in units of Li/ $\text{LiCoO}_2$  and (mA h)  $\text{g}^{-1}$  were determined from total charge transferred along with the film thickness and area assuming a film density equal to 85% of the bulk solid density as determined from ICP and thickness measurements.

## 3. Results and discussion

Typical film thicknesses were 100–900 nm for  $\text{LiCoO}_2$  films and 50–800 nm for  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  films. For the  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  films, a laser flux roughly 80 mJ pulse $^{-1}$  higher than that used for  $\text{LiCoO}_2$  films is required to achieve a similar deposition rate. ICP analysis shows a Li:Co ratio of  $1 \pm 0.05$  for the  $\text{LiCoO}_2$  films but Al leaching from the 7059 glass substrates has prevented an accurate determination of the metals stoichiometry in the aluminum containing films. However, the target has a 1:1 ratio of Co:Al and in PLD film growth the stoichiometry of the film is usually close to that of the target.

Fig. 1 shows X-ray diffraction  $\theta/2\theta$  spectra for a  $\text{LiCoO}_2$  film grown at  $T_s = 600^\circ\text{C}$ ,  $p[\text{O}_2] = 2000$  mTorr (curve a), for a  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  film grown at  $T_s = 600^\circ\text{C}$ ,  $p[\text{O}_2] = 1000$  mTorr (curve b) and for a bare  $\text{SnO}_2$  substrate (curve c). Spectra (a) and (b) have been vertically offset for clarity. The bottom panel shows, on a linear scale, the expected  $\text{LiCoO}_2$  powder diffraction pattern intensities and positions indexed to a hexagonal unit cell. Comparing the spectra for either the  $\text{LiCoO}_2$  or  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  sample with that of the bare  $\text{SnO}_2$  substrate clearly shows that most of the observed peaks are due to the substrate. However, for both films, the (003)  $\text{LiCoO}_2$  peak at  $\sim 19^\circ$  is the dominant peak with a scattering intensity comparable to that for the  $\text{SnO}_2$  substrate. The (104)  $\text{LiCoO}_2$  peak at  $\sim 45^\circ$ , which would be of comparable strength in a powder sample, is roughly 100 times weaker than the (003) peak for both the  $\text{LiCoO}_2$  and the  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  films. These spectra show that the films are highly (001) textured with the Li and Co layers

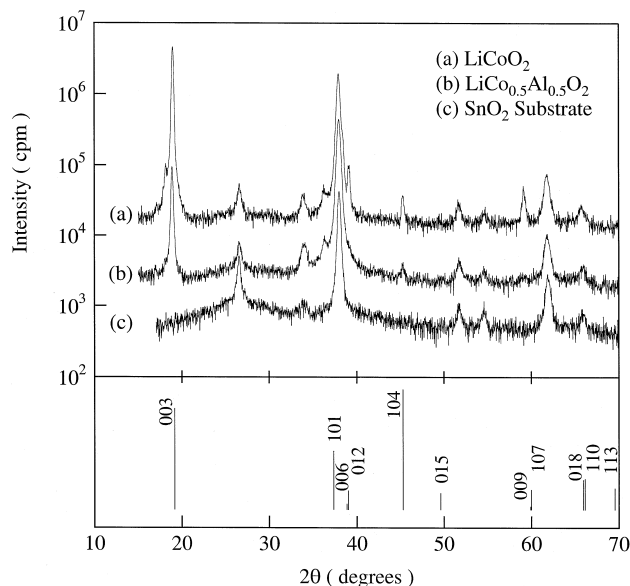


Fig. 1. Top panel: X-ray diffraction spectra for  $\text{LiCoO}_2$  (a) and  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  (b) thin films grown on  $\text{SnO}_2$ -coated 7059 glass substrates (c). The spectra are offset vertically for clarity. Bottom panel:  $\text{LiCoO}_2$  powder diffraction pattern positions and intensities from JCPDS card 44-145. Vertical scale for bottom panel is linear.

parallel to the substrate surface. Pole figures (not shown) show the textured grains to have randomly oriented  $a$ -axes in the plane parallel to the substrate.

Fig. 2 shows SEM micrographs of a  $\text{LiCoO}_2$  film (top) grown at  $T_s = 600^\circ\text{C}$ ,  $p[\text{O}_2] = 2000$  mTorr and a  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  film (bottom) grown at  $T_s = 600^\circ\text{C}$ ,  $p[\text{O}_2] = 1000$  mTorr. The  $\text{LiCoO}_2$  film has a grain size of  $\sim 100$  nm whereas the  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  film has substantially smaller grains, roughly 20–40 nm. Increasing the substrate temperature to  $T_s = 700^\circ\text{C}$  increases the grain size to 200–300 nm in  $\text{LiCoO}_2$  and 50–100 nm in  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$ .

Fig. 3 shows Raman scattering spectra for three representative  $\text{LiCoO}_2$  films as well as for a  $\text{Co}_3\text{O}_4$  and a  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  film. The intensity scale for all three  $\text{LiCoO}_2$  films and the  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  is the same but that for the  $\text{Co}_3\text{O}_4$  film has been reduced by a factor of 2.4 for clarity. The top spectra for a  $\text{LiCoO}_2$  film grown at  $T_s = 600^\circ\text{C}$ ,  $p[\text{O}_2] = 2000$  mTorr shows two well defined peaks at Raman shifts of  $\sim 486$   $\text{cm}^{-1}$  and  $\sim 596$   $\text{cm}^{-1}$  as expected for layered  $\text{LiCoO}_2$  [15–17]. Whereas four peaks, inconsistent with those observed, should have been observed for  $\text{LiCoO}_2$  with the cubic lithiated spinel structure [15]. In contrast, the X-ray diffraction spectra for the layered and lithiated-spinel phase of  $\text{LiCoO}_2$  differ only slightly. However, taken together, the measured X-ray diffraction and Raman scattering spectra unambiguously show that the films have the layered  $\text{LiCoO}_2$  structure. The spectra for the  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  film also shows the two-line layered-structure pattern with the peaks shifted slightly up in energy to 490 and 602  $\text{cm}^{-1}$ , respectively.

However, the shifts are small relative to the peak width for the  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  spectra.

The spectra for the  $\text{LiCoO}_2$  films grown at  $T_s = 400^\circ\text{C}$ ,  $p[\text{O}_2] = 200$  mTorr and  $T_s = 400^\circ\text{C}$ ,  $p[\text{O}_2] = 100$  mTorr show additional peaks at  $\sim 690$   $\text{cm}^{-1}$  and  $520$   $\text{cm}^{-1}$  with an additional weak but discernable peak at  $\sim 190$   $\text{cm}^{-1}$  in the latter film as well. Comparison of these spectra with that shown for the  $\text{Co}_3\text{O}_4$  film suggests that these extra peaks may arise due to a trace  $\text{Co}_3\text{O}_4$  impurity phase. However, no discernable  $\text{Co}_3\text{O}_4$  peaks are apparent in the X-ray diffraction spectra. For  $\text{LiCoO}_2$  films, the inset shows the strength of the  $690$   $\text{cm}^{-1}$  impurity peak relative to the  $596$   $\text{cm}^{-1}$   $\text{LiCoO}_2$  peak for  $200^\circ\text{C} < T_s < 800^\circ\text{C}$  ( $y$ -axis) and  $10$  mTorr  $< p[\text{O}_2] < 10,000$  mTorr ( $x$ -axis). An open circle, marking each point, is filled with a solid

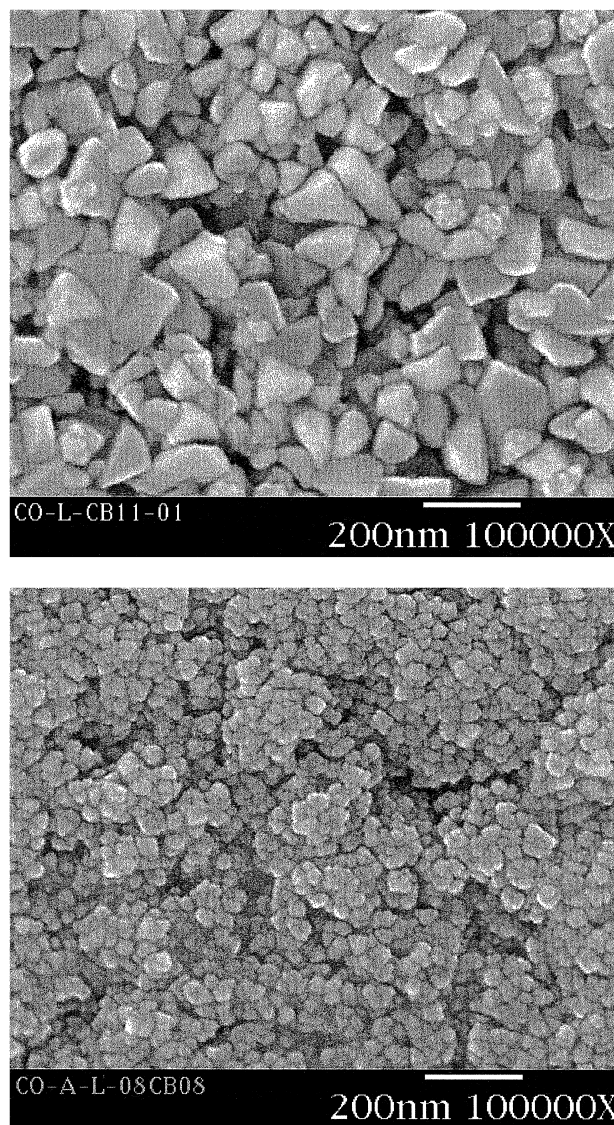


Fig. 2. Top: SEM micrograph of a  $\text{LiCoO}_2$  film grown at  $600^\circ\text{C}$  in 2 Torr  $\text{O}_2$  on  $\text{SnO}_2$ -coated glass. Bottom: SEM micrograph of a  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  film grown at  $600^\circ\text{C}$  in 1 Torr  $\text{O}_2$  on  $\text{SnO}_2$ -coated glass.

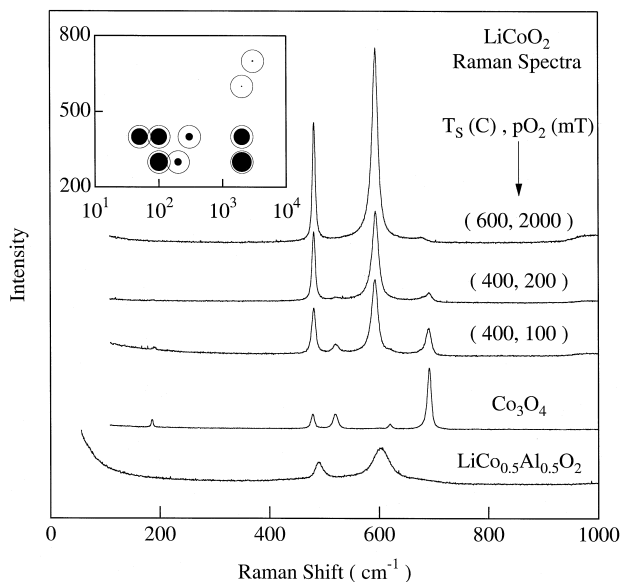


Fig. 3. Unpolarized Raman scattering spectra measured at 300 K for three representative  $\text{LiCoO}_2$  thin films as well as for a  $\text{Co}_3\text{O}_4$  and a  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  thin film. For each  $\text{LiCoO}_2$  spectra, the substrate temperature and oxygen partial pressure during film growth are shown in the figure. For  $\text{LiCoO}_2$  films, the inset shows the strength of the  $690\text{ cm}^{-1}$  impurity peak relative to the  $596\text{ cm}^{-1}$   $\text{LiCoO}_2$  peak for  $200^\circ\text{C} < T_s < 800^\circ\text{C}$  (y-axis) and  $10\text{ mTorr} < p[\text{O}_2] < 10,000\text{ mTorr}$  (x-axis). An open circle, marking each point, is filled with a solid circle whose diameter is proportional to  $I[690]/I[596]$ . For a fully filled circle,  $I[690]/I[596] = 0.5$ .

circle whose diameter is proportional to  $I[690]/I[596]$ . For a fully filled circle,  $I[690]/I[596] = 0.5$ . In general, the intensity of peak at  $690\text{ cm}^{-1}$  relative to that for the  $\text{LiCoO}_2$  peak at  $596\text{ cm}^{-1}$  decreases with increasing  $p[\text{O}_2]$  and  $T_s$  suggesting a reduced impurity content in the films grown at higher oxygen partial pressures and higher temperatures. However, for a fixed substrate temperature,  $400^\circ\text{C}$  for example, the relative impurity Raman scattering strength increases for  $p[\text{O}_2] > 1000\text{ mTorr}$ . This is due to a marked decrease in the strength of intrinsic  $\text{LiCoO}_2$  Raman scattering as opposed to an increase in the impurity phase scattering. This suggests that for a given substrate temperature, the optimum oxygen pressure is determined by balancing impurity phase reduction with  $\text{LiCoO}_2$  crystallinity.

The constant current charge and discharge capacities as a function of cycle number are shown in Fig. 4 for a  $\text{LiCoO}_2$  film grown at  $T_s = 600^\circ\text{C}$ ,  $p[\text{O}_2] = 200\text{ mTorr}$ . The film was cycled at  $5\text{ }\mu\text{A}$  between 3.8 and 4.2 V in a 1 M  $\text{LiClO}_4/\text{PPC}$  electrolyte. An initial discharge capacity of  $0.33\text{ Li/LiCoO}_2$  ( $89\text{ (mA h) g}^{-1}$ ) decreases to  $0.18\text{ Li/LiCoO}_2$  ( $49\text{ (mA h) g}^{-1}$ ) after 100 cycles. After 100 cycles, the rate of subsequent capacity loss is 0.25% per cycle. The inset shows the voltage vs. Li content for the first charge and discharge cycles. The shape of the charge and discharge curves are very similar, but only about 70% of the extracted Li can be reintercalated during discharge.

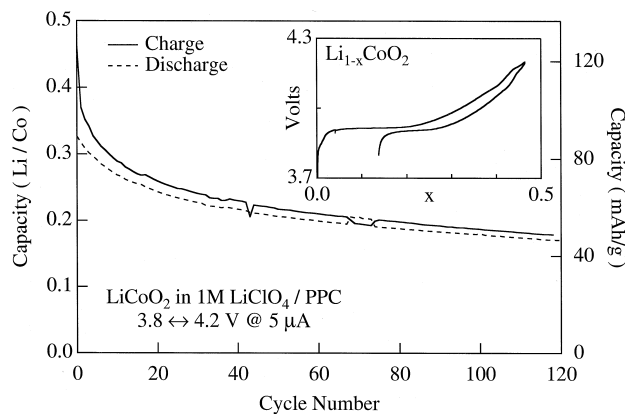


Fig. 4. Capacity, in units of Li exchanged per Co, vs. cycle number for a  $3900\text{ \AA}$  thick  $\text{LiCoO}_2$  film grown at  $600^\circ\text{C}$  in  $200\text{ mTorr}$  of oxygen. The film was cycled between 3.8 and 4.2 V vs. Li at constant current of  $5\text{ }\mu\text{A}$  in a 1 M  $\text{LiClO}_4/\text{PPC}$  electrolyte solution. The inset shows the voltage vs. Li content for the first charge and discharge cycles.

Fig. 5 shows initial electrochemical cycling results for a  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  film grown at  $T_s = 300^\circ\text{C}$ ,  $p[\text{O}_2] = 100\text{ mTorr}$ . The film was cycled at a constant current in 1 M  $\text{LiClO}_4/\text{PPC}$ . For the first 28 cycles, the film was cycled between 3.8 and 4.2 V at  $5\text{ }\mu\text{A}$ . While the charge capacity of the very first charge cycle is comparable to that for a  $\text{LiCoO}_2$  film, the first discharge capacity is substantially lower, about  $0.1\text{ Li/LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  ( $33\text{ (mA h) g}^{-1}$ ) compared to  $0.3\text{ Li/LiCoO}_2$  ( $82\text{ (mA h) g}^{-1}$ ) as shown in Fig. 4. After 20 cycles, the cell was left in an open circuit mode for an extended time compared to the typical cycling time. The charge capacity of the next charge cycle is roughly twice that of the previous charge cycle but the subsequent discharge capacity is essentially unchanged from the previous cycle. The inset shows the voltage vs. Li content for the first two charge and discharge cycles. In comparison with  $\text{LiCoO}_2$  (Fig. 4, inset), for  $\text{LiCo}_{0.5}$ -

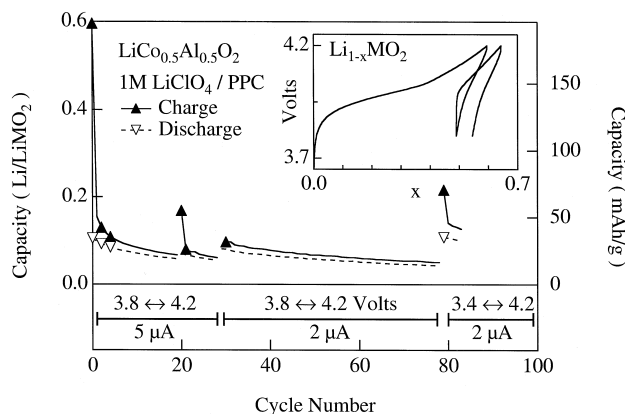


Fig. 5. Capacity, in units of Li exchanged per  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$ , vs. cycle number for a  $1700\text{ \AA}$  thick  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  film grown at  $300^\circ\text{C}$  in  $100\text{ mTorr}$  of oxygen. The cell was cycled at constant current in 1 M  $\text{LiClO}_4/\text{PPC}$ . During cycling the voltage limits and current were changed as indicated in the figure. The inset shows the voltage vs. Li content for the first two charge and discharge cycles.

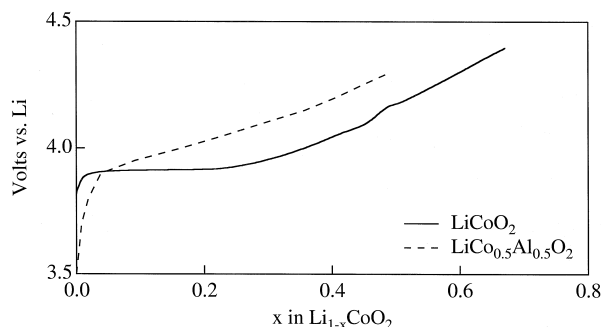


Fig. 6. Voltage vs. Li extracted on the first charge cycle for a  $\text{LiCoO}_2$  film (solid line) and a  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  film (dashed line). The  $3500 \text{ \AA}$  thick  $\text{LiCoO}_2$  film was grown at  $T_s = 700^\circ\text{C}$  and  $p[\text{O}_2] = 2000 \text{ mTorr}$  while the  $8100 \text{ \AA}$  thick  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  film was grown at  $T_s = 600^\circ\text{C}$  and  $p[\text{O}_2] = 1000 \text{ mTorr}$ .

$\text{Al}_{0.5}\text{O}_2$ , the discharge does not follow the charge curve and only about 18% of the extracted Li can be reintercalated. The cycling behavior  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  along with the large polarization of the voltage vs. charge curves relative to that for the  $\text{LiCoO}_2$  films suggests a substantial asymmetry between Li extraction and intercalation dynamics in  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$ . After 28 cycles, the current was reduced from 5 to 2  $\mu\text{A}$  while keeping the voltage range the same as indicated in the bottom panel of Fig. 5. This factor of 2.5 reduction in the current increased the discharge capacity 40%, but the discharge capacity is still quite low compared to either  $\text{LiCoO}_2$  or the first-cycle charge capacity of this same  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  film. After 80 cycles the discharge voltage limit was reduced from 3.8 to 3.4 V which resulted in a one cycle increase in capacity of 50% but with a increased rate of capacity loss per cycle as well.

Fig. 6 compares the voltage as a function of Li extracted during the first charge half-cycle for a  $\text{LiCoO}_2$  film (solid line) and a  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  film (dashed line). During most of the charge cycle, the voltage of  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  cathode is higher than that of the  $\text{LiCoO}_2$  cathode in agreement with expectations from recent first-principles calculations [10]. However, as clear from Fig. 5, for the  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  films grown to date, neither an increased voltage nor a reasonable capacity are realized upon discharge. We also note that the voltage plateau indicative of a two phase region for  $x < 0.2$  in  $\text{Li}_{1-x}\text{CoO}_2$  is not observed for  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$ .

#### 4. Summary

$\text{LiCoO}_2$  and  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  thin films have been grown by pulsed laser ablation on  $\text{SnO}_2$  coated glass substrates.

For both stoichiometries, the resultant films are dense and uniaxially textured with the Li and Co layers parallel to the substrate.  $\text{LiCoO}_2$  films grown at  $T_s = 600^\circ\text{C}$ ,  $p[\text{O}_2] = 2000 \text{ mTorr}$  have a grain size of  $\sim 100 \text{ nm}$  and reasonable electrochemical properties. The  $\text{LiCo}_{0.5}\text{Al}_{0.5}\text{O}_2$  films, on the other hand, have roughly 3 times less capacity than the  $\text{LiCoO}_2$  films and apparently a large asymmetry between Li extraction and reintercalation. Future film growth and characterization work will focus on elucidating the cause of, and hopefully reducing, the present decreased capacity in the aluminum substituted films.

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#### References

- [1] K. Mizushima, P.C. Jones, P.J. Wiseman, J.B. Goodenough, *Mater. Res. Bull.* 15 (1980) 783.
- [2] T.A. Hewston, B.L. Chamberland, *J. Phys. Chem. Solids* 48 (1987) 97.
- [3] K. Ozawa, *Solid State Ionics* 69 (1994) 212.
- [4] C. Delmas, H. Cognac-Auradou, J.M. Cocciantelli, M. Menetrier, J.P. Doumerc, *Solid State Ionics, Diffusion and Reactions* 69 (1994) 257.
- [5] G.G. Amatucci, J.M. Tarascon, L.C. Klein, *JECS* 143 (1996) 1114.
- [6] C. Wolverton, A. Zunger, *JECS* 145 (1998) 2424.
- [7] J.N. Reimers, J.R. Dahn, U. von Sacken, *JECS* 140 (1993) 2752.
- [8] G.G. Amatucci, J.M. Tarascon, L.C. Klein, *Solid State Ionics* 83 (1996) 167.
- [9] Y.-M. Choi, S.-I. Pyun, J.-S. Bae, S.-I. Moon, *J. Power Sources* 56 (1995) 25.
- [10] G. Ceder, Y.-M. Chiang, D.R. Sadoway, M.K. Aydinol, Y.-I. Jang, B. Huang, *Nature* 392 (1998) 694.
- [11] M. Antaya, K. Cearn, J.S. Preston, J.N. Reimers, J.R. Dahn, *J. Appl. Phys.* 76 (1994) 2799.
- [12] K.A. Striebel, C.Z. Deng, S.J. Wen, E.J. Cairns, *JECS* 143 (1996) 1821.
- [13] R.J. Gummow, D.C. Liles, M.M. Thackeray, *Mater. Res. Bull.* 28 (1993) 235.
- [14] E. Rossen, J.N. Reimers, J.R. Dahn, *Solid State Ionics* 62 (1993) 53.
- [15] W. Huang, R. Frech, *Solid State Ionics* 86–88 (1996) 395.
- [16] M. Inaba, Y. Iriyama, Z. Ogumi, Y. Todzuka, A. Tasaka, *J. Raman Spectrosc.* 28 (1997) 613.
- [17] J.D. Perkins, M.L. Fu, D.M. Trickett, J.M. McGraw, T.F. Ciszek, P.A. Parilla, C.T. Rogers, D.S. Ginley, *Mater. Res. Soc. Symp. Proc.* 496 (1998) 329.